Revisiting the quasinormal modes of the Schwarzschild black hole: Numerical analysis<br>Luis Alex Huahuachampi Mamani (UFRB)<br>in collaboration with<br>Lucas Sanches, Vilson Zanchin (UFABC), and Angel Masa (U. de Cartagena) Eur.Phys.J.C 82 (2022) 10, 897<br>VII Join Meeting of Graduate Students Physics - Física/UNSA the 10th of November, 2022<br>Arequipa - Perú

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## Einstein equations

■ Einstein's equations proposed em 1915 are

$$
\begin{equation*}
G_{\mu \nu}=8 \pi G T_{\mu \nu} \tag{1}
\end{equation*}
$$

- Schwarzschild' solution: spherically solution found by Karl Schwarzschild in 1916. The metric is given by

$$
\begin{equation*}
d s^{2}=-f(r) d t^{2}+\frac{1}{f(r)} d r^{2}+r^{2} d \theta^{2}+r^{2} \sin ^{2} \theta d \phi^{2}, \quad \text { onde } \quad f(r)=1-\frac{2 M}{r^{2}} . \tag{2}
\end{equation*}
$$

- Perturbation equations: the radial equation can be written in the Schrödinger-like form

$$
\begin{equation*}
\frac{d^{2} \psi_{s}(r)}{d r_{*}^{2}}+\left(\omega^{2}-V_{s}(r)\right) \psi_{s}(r)=0 \tag{3}
\end{equation*}
$$

## Perturbation equations - Bosonic fields

The potential of the Schrödinger-like equation for bosonic fields is given by (see for instance [Berti-Cardoso-Starinets, 2009])

$$
\begin{equation*}
V_{s}(r)=f(r)\left(\frac{\ell(\ell+1)}{r^{2}}+\left(1-s^{2}\right) \frac{2 M}{r^{3}}\right) \tag{4}
\end{equation*}
$$

After applying a few transformations the differential equation suitable to apply the pseudo-spectral and AIM methods is

$$
\begin{align*}
\left(u\left(\ell(\ell+1)-s^{2} u\right)-4\right. & \left.i \lambda-16 u(1+u) \lambda^{2}\right) \phi_{s}(u)  \tag{5}\\
& +\left(u^{3}+4 i u\left(1-2 u^{2}\right) \lambda\right) \phi_{s}^{\prime}(u)-(1-u) u^{3} \phi_{s}^{\prime \prime}(u)=0
\end{align*}
$$

where $r=1 / u$ and $\lambda=\omega M$. Note that $u \in[0,1]$.

## Perturbation equation - spin $1 / 2$ field

For spin $1 / 2$ field the potential of the Schrödinger-like equation is given by (see [Cho, Phys. Rev. D 68, 024003] and Shu and Shen, Phys. Lett. B 619, 340])

$$
\begin{equation*}
V_{1 / 2}=\frac{(1+\ell) \sqrt{r-2 M}}{r^{7 / 2}}((1+\ell) \sqrt{r(r-2 M)}+3 M-r) \tag{6}
\end{equation*}
$$

After applying a few transformation the differential equation suitable to apply the pseudospectral and AIM methods is

$$
\begin{aligned}
& \left(u^{3}+u(1+\ell)(1+\ell-\sqrt{1-u})+\frac{u^{2}}{2}\left((1+\ell)(3 \sqrt{1-u}-4)-2 \ell^{2}\right)-4 i(1-u) \lambda\right. \\
& \left.-16 u\left(1-u^{2}\right) \lambda^{2}\right) \phi_{1 / 2}(u)+\left(u^{3}(1-u)+4 i u\left(1-u-2 u^{2}+2 u^{3}\right) \lambda\right) \phi_{1 / 2}^{\prime}(u) \\
& -u^{3}(1-u)^{2} \phi_{1 / 2}^{\prime \prime}(u)=0
\end{aligned}
$$

where $r=1 / u$ and $\lambda=\omega M$. Note that $u \in[0,1]$. For spin $3 / 2$ and $5 / 2$ fields we did the same analysis, we did not present here because their expressions are huge.

## The pseudo-spectral method - short review

The idea behind the pseudo-spectral method is to rewrite the regular function $\phi(u)$ in a base composed by cardinal functions $C_{j}(u)$ and the Gauss-Lobato grid, in the form (for details see [Jansen, Eur. Phys. J. Plus 132, 546])

$$
\begin{equation*}
\phi_{s}(u)=\sum_{j=0}^{N} g\left(u_{j}\right) C_{j}(u) ; \quad u_{i}=\frac{1}{2}\left(1 \pm \cos \left[\frac{i}{N} \pi\right]\right), \quad i=0,1,2, \cdots, N \tag{7}
\end{equation*}
$$

Then, the matrix representation of the quadratic eigenvalue problem can be written as

$$
\begin{equation*}
\left(\tilde{M}_{0}+\tilde{M}_{1} \lambda+\tilde{M}_{2} \lambda^{2}\right) g=0 ; \quad \rightarrow \quad\left(M_{0}+M_{1} \lambda\right) \cdot \vec{g}=\mathbb{0} \tag{8}
\end{equation*}
$$

$$
C_{j}(u)=T_{j}(u), \quad C_{j}(u)=\frac{2}{N p_{j}} \sum_{m=0}^{N} \frac{1}{p_{m}} T_{m}\left(u_{j}\right) T_{m}(u), \quad\left\{\begin{array}{l}
p_{0}=2  \tag{9}\\
p_{N}=2 \\
p_{j}=1
\end{array}\right.
$$

## The asymptotic iteration method - short review

The second order differential equation can be written in the form [H. Ciftci et al,. 2003, Journal of Physics A: Mathematical and General]

$$
\begin{equation*}
y^{\prime \prime}(x)-\lambda_{0}(x) y^{\prime}(x)-s_{0}(x) y(x)=0 \tag{10}
\end{equation*}
$$

where $\lambda_{0}(x) \neq 0$ and $s_{0}(x)$ are $C_{\infty}$. A general solution of Eq. (10) can be written in the form

$$
\begin{equation*}
y(x)=\exp \left(-\int \alpha \mathrm{d} t\right) \times\left[C_{2}+C_{1} \int^{x} \exp \left(\int^{t}\left(\lambda_{0}(\tau)+2 \alpha(\tau)\right) \mathrm{d} \tau\right) \mathrm{d} t\right] \tag{11}
\end{equation*}
$$

if for some $n>0$ the condition

$$
\begin{equation*}
\delta \equiv s_{n} \lambda_{n-1}-\lambda_{n} s_{n-1}=0 \tag{12}
\end{equation*}
$$

is satisfied. However, for computing QNMs it is necessary a modification of this procedure in order to circumvent numerical problems.

## The asymptotic iteration method - short review

Instead, the coefficients are expanded around an arbitrary point, say $x=\xi$, (see [Cho, Cornell, Doukas, Huang, and Naylor, Adv. Math. Phys.2012, 281705])

$$
\begin{equation*}
\lambda_{n}(\xi)=\sum_{i=0}^{\infty} c_{n}^{i}(x-\xi)^{i} ; \quad s_{n}(\xi)=\sum_{i=0}^{\infty} d_{n}^{i}(x-\xi)^{i} \tag{13}
\end{equation*}
$$

The quantization condition becomes

$$
\begin{equation*}
\delta \equiv d_{n}^{0} c_{n-1}^{0}-d_{n-1}^{0} c_{n}^{0}=0 \tag{14}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{n}^{i}=(i+1) c_{n-1}^{i+1}+d_{n-1}^{i}+\sum_{k=0}^{i} c_{0}^{k} c_{n-1}^{i-k} ; \quad d_{n}^{i}=(i+1) d_{n-1}^{i+1}+\sum_{k=0}^{i} d_{0}^{k} c_{n-1}^{i-k} \tag{15}
\end{equation*}
$$

A package implemented in Julia can be found in: https://github.com/lucass-carneiro/QuasinormalModes.jl.

## Spin 0

| $l$ | $n$ | Pseudo-spectral <br> I (60 Polynomials) | Pseudo-spectral <br> II (40 polynomials) | AIM <br> 100 Iterations | Ref. [16] | Ref. [17] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | $\pm 0.110455-0.104896 i$ | $\pm 0.110455-0.104896 i$ | $\pm 0.110455-0.104896 i$ | $0.1046-0.1152 i$ | $\pm 0.1105-0.1008 i$ |
| 1 | 0 | $\pm 0.292936-0.097660 i$ | $\pm 0.292936-0.097660 i$ | $\pm 0.292936-0.097660 i$ | $0.2911-0.0980 i$ | $\pm 0.2929-0.0978 i$ |
|  | 1 | $\pm 0.264449-0.306257 i$ | $\pm 0.264449-0.306257 i$ | - | - | $\pm 0.2645-0.3065 i$ |
| 2 | 0 | $\pm 0.483644-0.096759 i$ | $\pm 0.483644-0.096759 i$ | $\pm 0.483644-0.096759 i$ | $0.4832-0.0968 i$ | $\pm 0.4836-0.0968 i$ |
|  | 1 | $\pm 0.463851-0.295604 i$ | $\pm 0.463851-0.295604 i$ | $\pm 0.463851-0.295604 i$ | $0.4632-0.2958 i$ | $\pm 0.4638-0.2956 i$ |
|  | 2 | $\pm 0.430544-0.508558 i$ | $\pm 0.430544-0.508558 i$ | - | - | $\pm 0.4304-0.5087 i$ |
| 3 | 0 | $\pm 0.675366-0.096500 i$ | $\pm 0.675366-0.096500 i$ | $\pm 0.675366-0.096500 i$ | $0.6752-0.0965 i$ | - |
|  | 1 | $\pm 0.660671-0.292285 i$ | $\pm 0.660671-0.292285 i$ | $\pm 0.660671-0.292285 i$ | $0.6604-0.2923 i$ | - |
|  | 2 | $\pm 0.633626-0.496008 i$ | $\pm 0.633626-0.496008 i$ | $\pm 0.633626-0.496008 i$ | $0.6348-0.4941 i$ | - |
|  | 3 | $\pm 0.598773-0.711221 i$ | $\pm 0.598769-0.711220 i$ | - | - | - |
| 4 | 0 | $\pm 0.867416-0.096392 i$ | $\pm 0.867416-0.096392 i$ | $\pm 0.867416-0.096392 i$ | $0.8673-0.0964 i$ | - |
|  | 1 | $\pm 0.855808-0.290876 i$ | $\pm 0.855808-0.290876 i$ | $\pm 0.855808-0.290876 i$ | $0.8557-0.2909 i$ | - |
|  | 2 | $\pm 0.833692-0.490325 i$ | $\pm 0.833692-0.490325 i$ | $\pm 0.833692-0.490325 i$ | $0.8345-0.4895 i$ | - |
|  | 3 | $\pm 0.803288-0.697482 i$ | $\pm 0.803288-0.697482 i$ | $\pm 0.803288-0.697482 i$ | $0.8064-0.6926 i$ | - |
|  | 4 | $\pm 0.767733-0.914019 i$ | $\pm 0.767679-0.914105 i$ | - | - | - |

[16] Shu and Shen, Phys. Lett. B 619, 340 (2005), [17] Konoplya, J. Phys. Stud. 8, 93 (2004).

## Spin 1

| $l$ | $n$ | Pseudo-spectral <br> I (60 Polynomials) | Pseudo-spectral <br> II (40 polynomials $)$ | AIM <br> 100 Iterations | Ref. [16] | Ref. [17] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | - | - | - | - | - |
| 1 | 0 | $\pm 0.248263-0.092488 i$ | $\pm 0.248263-0.092488 i$ | $\pm 0.248263-0.092488 i$ | $0.2459-0.0931 i$ | $\pm 0.2482-0.0926 i$ |
|  | 1 | $\pm 0.214515-0.293668 i$ | $\pm 0.214515-0.293667 i$ | - | - | $\pm 0.2143-0.2941 i$ |
| 2 | 0 | $\pm 0.457596-0.095004 i$ | $\pm 0.457595-0.095004 i$ | $\pm 0.457596-0.095004 i$ | $0.4571-0.0951 i$ | $\pm 0.4576-0.0950 i$ |
|  | 1 | $\pm 0.436542-0.290710 i$ | $\pm 0.436542-0.290710 i$ | $\pm 0.436542-0.290710 i$ | $0.4358-0.2910 i$ | $\pm 0.4365-0.2907 i$ |
|  | 2 | $\pm 0.401187-0.501587 i$ | $\pm 0.401187-0.501587 i$ | - | - | $\pm 0.4009-0.5017 i$ |
| 3 | 0 | $\pm 0.656899-0.095616 i$ | $\pm 0.656899-0.095616 i$ | $\pm 0.656899-0.095616 i$ | $0.6567-0.0956 i$ | $\pm 0.6569-0.0956 i$ |
|  | 1 | $\pm 0.641737-0.289728 i$ | $\pm 0.641737-0.289728 i$ | $\pm 0.641737-0.289728 i$ | $0.6415-0.2898 i$ | $\pm 0.6417-0.2897 i$ |
|  | 2 | $\pm 0.613832-0.492066 i$ | $\pm 0.613832-0.492066 i$ | $\pm 0.613832-0.492066 i$ | $0.6151-0.4901 i$ | $\pm 0.6138-0.4921 i$ |
|  | 3 | $\pm 0.577919-0.706331 i$ | $\pm 0.577915-0.706328 i$ | - | - | $\pm 0.5775-0.7065 i$ |
| 4 | 0 | $\pm 0.853095-0.095860 i$ | $\pm 0.853095-0.095860 i$ | $\pm 0.853095-0.095810 i$ | $0.8530-0.0959 i$ | - |
|  | 1 | $\pm 0.841267-0.289315 i$ | $\pm 0.841267-0.289315 i$ | $\pm 0.841267-0.289315 i$ | $0.8411-0.2893 i$ | - |
|  | 2 | $\pm 0.818728-0.487838 i$ | $\pm 0.818728-0.487838 i$ | $\pm 0.818728-0.487838 i$ | $0.8196-0.4870 i$ | - |
|  | 3 | $\pm 0.787748-0.694242 i$ | $\pm 0.787748-0.694243 i$ | $\pm 0.787748-0.694242 i$ | $0.7909-0.6892 i$ | - |
| 4 | $\pm 0.751549-0.910242 i$ | $\pm 0.751481-0.910301 i$ | - | - | - |  |

[16] Shu and Shen, Phys. Lett. B 619, 340 (2005), [17] Konoplya, J. Phys. Stud. 8, 93 (2004).

## Spin 2

| $l$ |  | Pseudo-spectral I (60 Polynomials) | $\begin{aligned} & \text { Pseudo-spectral } \\ & \text { II (40 polynomials) } \end{aligned}$ | AIM <br> 100 Iterations | Ref. [16] | Ref. [17] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 |  | - | - | - | - |  |
| 1 | 0 1 | $\pm 0.110455-0.104896 i$ | $\pm 0.110455-0.104896 i$ - |  |  |  |
| 2 | 0 1 2 | $\pm 0.373672-0.088962 i$ $\pm 0.346711-0.273915 i$ $\pm 0.301053-0.478277 i$ | $\begin{aligned} & \pm 0.373672-0.088962 i \\ & \pm 0.346711-0.273915 i \\ & \pm 0.301053-0.478277 i \end{aligned}$ | $\begin{gathered} \pm 0.373672-0.088962 i \\ \pm 0.346711-0.273915 i \\ - \end{gathered}$ | $\begin{gathered} 0.3730-0.0891 i \\ 0.3452-0.2746 i \\ - \end{gathered}$ | $\begin{aligned} & \pm 0.3736-0.0890 i \\ & \pm 0.3463-0.2735 i \\ & \pm 0.2985-0.4776 i \end{aligned}$ |
| 3 | 0 1 2 3 | $\begin{aligned} & \pm 0.599443-0.092703 i \\ & \pm 0.582644-0.281298 i \\ & \pm 0.551685-0.479093 i \\ & \pm 0.511962-0.690337 i \end{aligned}$ | $\begin{aligned} & \pm 0.599443-0.092703 i \\ & \pm 0.582644-0.281298 i \\ & \pm 0.551685-0.479093 i \\ & \pm 0.511966-0.690333 i \end{aligned}$ | $\begin{gathered} \pm 0.599443-0.092703 i \\ \pm 0.582644-0.281298 i \\ \pm 0.551685-0.479027 i \\ - \\ \hline \end{gathered}$ | $0.5993-0.0927 i$ $0.5824-0.2814 i$ $0.5532-0.4767 i$ - | $\begin{aligned} & \pm 0.5994-0.0927 i \\ & \pm 0.5826-0.2813 i \\ & \pm 0.5516-0.4790 i \\ & \pm 0.5111-0.6905 i \\ & \hline \end{aligned}$ |
| 4 | 2 <br> 2 <br> 3 <br> 4 | $\pm 0.809178-0.094164 i$ $\pm 0.796632-0.284334 i$ $\pm 0.772710-0.479908 i$ $\pm 0.739837-0.683924 i$ $\pm 0.701516-0.898239 i$ | $\pm 0.809178-0.094164 i$ <br> $\pm 0.796632-0.284334 i$ <br> $\pm 0.772710-0.479908 i$ <br> $\pm 0.739836-0.683925 i$ <br> $\pm 0.701398-0.898196 i$ | $\pm 0.809178-0.094164 i$ <br> $\pm 0.796632-0.284334 i$ <br> $\pm 0.772710-0.479908 i$ <br> $\pm 0.739837-0.683924 i$ <br> - | $0.8091-0.0942 i$ $0.7965-0.2844 i$ $0.7736-0.4790 i$ $0.7433-0.6783 i$ - | $\pm 0.8092-0.0942 i$ $\pm 0.7966-0.2843 i$ $\pm 0.7727-0.4799 i$ $\pm 0.7397-0.6839 i$ $\pm 0.7006-0.8985 i$ |

[16] Shu and Shen, Phys. Lett. B 619, 340 (2005), [17] Konoplya, J. Phys. Stud. 8, 93 (2004).

## Spin 1/2

| $l$ | $n$ | Pseudo-spectral <br> I $(60$ Polynomaials $)$ | Pseudo-spectral <br> II (40 polynomials) | AIM <br> 100 Iterations | Ref. [16] | Ref. [18] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | $\pm 0.182963-0.096982 i$ | $\pm 0.182963-0.096982 i$ | $\pm 0.182963-0.096824 i$ | - | - |
| 1 | 0 | $\pm 0.380037-0.096405 i$ | $\pm 0.380037-0.096405 i$ | $\pm 0.380037-0.096405 i$ | $0.3786-0.0965 i$ | $0.379-0.097 i$ |
|  | 1 | $\pm 0.355833-0.297497 i$ | $\pm 0.355833-0.297497 i$ | $\pm 0.355833-0.297497 i$ | - |  |
| 2 | 0 | $\pm 0.574094-0.096305 i$ | $\pm 0.574094-0.096305 i$ | $\pm 0.574094-0.096305 i$ | $0.5737-0.0963 i$ | $0.574-0.096 i$ |
|  | 1 | $\pm 0.557015-0.292715 i$ | $\pm 0.557015-0.292715 i$ | $\pm 0.557015-0.292715 i$ | $0.5562-0.2930 i$ | $0.556-0.293 i$ |
|  | 2 | $\pm 0.526607-0.499695 i$ | $\pm 0.526607-0.499695 i$ | $\pm 0.526607-0.499695 i$ | - | - |
| 3 | 0 | $\pm 0.767355-0.096270 i$ | $\pm 0.767355-0.096270 i$ | $\pm 0.767355-0.096270 i$ | $0.7672-0.0963 i$ | $0.767-0.096 i$ |
|  | 1 | $\pm 0.754300-0.290968 i$ | $\pm 0.754300-0.290968 i$ | $\pm 0.754300-0.290968 i$ | $0.7540-0.2910 i$ | $0.754-0.291 i$ |
|  | 2 | $\pm 0.729770-0.491910 i$ | $\pm 0.729770-0.491910 i$ | $\pm 0.729770-0.491910 i$ | $0.7304-0.4909 i$ | $0.730-0.491 i$ |
|  | 3 | $\pm 0.696913-0.702293 i$ | $\pm 0.696913-0.702293 i$ | $\pm 0.696913-0.702293 i$ | - | - |
| 4 | 0 | $\pm 0.960293-0.096254 i$ | $\pm 0.960292-0.096254 i$ | $\pm 0.960293-0.096254 i$ | $0.9602-0.0963 i$ | $0.960-0.096 i$ |
|  | 1 | $\pm 0.949759-0.290148 i$ | $\pm 0.949759-0.290148 i$ | $\pm 0.949759-0.290148 i$ | $0.9496-0.2902 i$ | $0.950-0.290 i$ |
|  | 2 | $\pm 0.929494-0.488116 i$ | $\pm 0.929494-0.488116 i$ | $\pm 0.929494-0.488116 i$ | $0.9300-0.4876 i$ | $0.930-0.488 i$ |
|  | 3 | $\pm 0.901129-0.692520 i$ | $\pm 0.901129-0.692520 i$ | $\pm 0.901129-0.692520 i$ | $0.9036-0.6892 i$ | $0.904-0.689 i$ |
|  | 4 | $\pm 0.867043-0.905047 i$ | $\pm 0.867008-0.905066 i$ | $\pm 0.867043-0.905047 i$ | - | - |

[16] Shu and Shen, Phys. Lett. B 619, 340 (2005), [18] Cho, Phys. Rev. D 68, 024003 (2003).

## Spin 3/2

| $l$ | $n$ | Pseudo-spectral <br> I (60 Polynomials) | Pseudo-spectral <br> II (40 polynomials) | AIM <br> 100 Iterations | Ref. [16] | Ref. [23] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | $\pm 0.311292-0.090087 i$ | $\pm 0.311292-0.090087 i$ | $\pm 0.311292-0.090087 i$ | - | $0.3112-0.0902 i$ |
| 1 | 0 | $\pm 0.530048-0.093751 i$ | $\pm 0.530048-0.093751 i$ | $\pm 0.530048-0.093751 i$ | - | $0.5300-0.0937 i$ |
|  | 1 | $\pm 0.511392-0.285423 i$ | $\pm 0.511392-0.285423 i$ | $\pm 0.511392-0.285423 i$ | - | $0.5113-0.2854 i$ |
| 2 | 0 | $\pm 0.734750-0.094878 i$ | $\pm 0.734750-0.094878 i$ | $\pm 0.734750-0.094878 i$ | $0.7346-0.0949 i$ | $0.7347-0.0948 i$ |
|  | 1 | $\pm 0.721047-0.286906 i$ | $\pm 0.721047-0.286906 i$ | $\pm 0.721047-0.286906 i$ | $0.7206-0.2870 i$ | $0.7210-0.2869 i$ |
|  | 2 | $\pm 0.695287-0.485524 i$ | $\pm 0.695287-0.485524 i$ | $\pm 0.695287-0.485524 i$ | - | $0.6952-0.4855 i$ |
| 3 | 0 | $\pm 0.934364-0.095376 i$ | $\pm 0.934364-0.095376 i$ | $\pm 0.934364-0.095376 i$ | $0.9343-0.0954 i$ | $0.9343-0.0953 i$ |
|  | 1 | $\pm 0.923502-0.287560 i$ | $\pm 0.923502-0.287560 i$ | $\pm 0.923502-0.287560 i$ | $0.9233-0.2876 i$ | $0.9235-0.2875 i$ |
|  | 2 | $\pm 0.902599-0.483957 i$ | $\pm 0.902599-0.483957 i$ | $\pm 0.902599-0.483957 i$ | $0.9031-0.4835 i$ | $0.9025-0.4839 i$ |
|  | 3 | $\pm 0.873342-0.687024 i$ | $\pm 0.873343-0.687024 i$ | $\pm 0.873342-0.687024 i$ | - | $0.8732-0.6870 i$ |
| 4 | 0 | $\pm 1.131530-0.095640 i$ | $\pm 1.131530-0.095640 i$ | $\pm 1.131530-0.095640 i$ | $1.1315-0.0956 i$ | $1.1315-0.0956 i$ |
|  | 1 | $\pm 1.122523-0.287908 i$ | $\pm 1.122523-0.287908 i$ | $\pm 1.122523-0.287908 i$ | $1.1224-0.2879 i$ | $1.1225-0.2879 i$ |
|  | 2 | $\pm 1.104976-0.483096 i$ | $\pm 1.104976-0.483096 i$ | $\pm 1.104976-0.483096 i$ | $1.1053-0.4828 i$ | $1.1049-0.4830 i$ |
|  | 3 | $\pm 1.079852-0.683000 i$ | $\pm 1.079852-0.683000 i$ | $\pm 1.079852-0.683000 i$ | $1.0817-0.6812 i$ | $1.0798-0.6829 i$ |
|  | 4 | $\pm 1.048599-0.889113 i$ | $\pm 1.048596-0.889115 i$ | $\pm 1.048599-0.889113 i$ | - | $1.0484-0.8890 i$ |

[16] Shu and Shen, Phys. Lett. B 619, 340 (2005), [23] Chen, Cho, Cornell, and Harmsen,

## Spin $1 / 2$ and $3 / 2$

It is worth mentioning that we found additional frequencies which are purely imaginary and the same for spin $1 / 2$ and $3 / 2$ fields.

| Pseudo-spectral | Pseudo-spectral | AIM |
| :---: | :---: | :---: |
| I (60 Polynomials) | II (40 polynomials) | 100 Iterations |
| -0.250000i | -0.250000i | -0.250000i |
| $-0.500000 i$ | -0.500000i | -0.500000i |
| -0.750000i | -0.750000i | -0.750000i |
| $-1.000000 i$ | -1.000000i | $-1.000031 i$ |
| -1.2499998i | $-1.250000 i$ | $-1.246550 i$ |

## Spin 5/2

| $l$ | $n$ | Pseudo-spectral <br> I $(60$ Polynomials $)$ | Pseudo-spectral <br> II $(40$ polynomials $)$ | AIM <br> 100 Iterations |
| :---: | :---: | :---: | :---: | :---: |
| 0 | 0 | $\pm 0.462727-0.092578 i$ | $\pm 0.462727-0.092578 i$ | $0.462727-0.092577 i$ |
| 1 | 0 | $\pm 0.687103-0.094566 i$ | $\pm 0.687103-0.094566 i$ | $0.687103-0.094566 i$ |
|  | 1 | $\pm 0.670542-0.285767 i$ | $\pm 0.670542-0.285767 i$ | $0.670542-0.285767 i$ |
| 2 | 0 | $\pm 0.897345-0.095309 i$ | $\pm 0.897345-0.095309 i$ | $0.897345-0.095309 i$ |
|  | 1 | $\pm 0.884980-0.287266 i$ | $\pm 0.884980-0.287266 i$ | $0.884980-0.287266 i$ |
|  | 2 | $\pm 0.861109-0.483113 i$ | $\pm 0.861109-0.483113 i$ | $0.861109-0.483113 i$ |
| 3 | 0 | $\pm 1.101190-0.095648 i$ | $\pm 1.101190-0.095648 i$ | $1.101190-0.095648 i$ |
|  | 1 | $\pm 1.091300-0.287886 i$ | $\pm 1.091300-0.287886 i$ | $1.091300-0.287886 i$ |
|  | 2 | $\pm 1.071999-0.482895 i$ | $\pm 1.071999-0.482895 i$ | $1.071999-0.482895 i$ |
|  | 3 | $\pm 1.044272-0.682307 i$ | $\pm 1.044272-0.682307 i$ | $1.044272-0.682307 i$ |
| 4 | 0 | $\pm 1.301587-0.095829 i$ | $\pm 1.301587-0.095829 i$ | $1.301587-0.095829 i$ |
|  | 1 | $\pm 1.293328-0.288184 i$ | $\pm 1.293328-0.288184 i$ | $1.293328-0.288184 i$ |
| 2 | $\pm 1.277107-0.482604 i$ | $\pm 1.277107-0.482604 i$ | $1.277107-0.482604 i$ |  |
|  | 3 | $\pm 1.253526-0.680366 i$ | $\pm 1.253526-0.680366 i$ | $1.253526-0.680366 i$ |
|  | $41.223513-0.882554 i$ | $\pm 1.223512-0.882553 i$ | $1.223513-0.882554 i$ |  |

## Spin 5/2

We also found purely imaginary frequencies for spin $5 / 2$ field.

| Pseudo-spectral <br> I (60 Polynomials) $)$ | Pseudo-spectral | AIM |
| :---: | :---: | :---: |
| $-0.125000 i$ | $-0.125000 i$ | $-0.125000 i$ |
| $-0.375602 i$ | $-0.375602 i$ | $-0.378659 i$ |
| $-0.626877 i$ | $-0.626877 i$ | $-0.623931 i$ |
| $-0.878946 i$ | $-0.878948 i$ | $-0.907374 i$ |

## Conclusions and outlook

■ We applied successfully the pseudo-spectral and AIM methods to calculate quasinormal frequencies of the Schwarzschild black hole for spin $0,1 / 2,1,3 / 2,2$ and 5/2 fields.
■ For bosonic fields (spin 0, 1, and 2) we found results in agreement with results available in the literature. We did not find any additional frequency.

■ For fermionic fields (spin $1 / 2$ and $3 / 2$ ) we found results in agreement with results available in the literature. We also found additional frequencies which are purely imaginary.

- For spin $5 / 2$ field we found results not available in the literature. We also found frequencies which are purely imaginary.
■ In the next stage we are going to address the case of charged and rotating black holes.


## Thank you!

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